# Vega's hot dust from icy planetesimals scattered inwards by an outward-migrating planetary system

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Accepted 2014 March 19. Received 2014 March 17; in original form 2014 February 5

### ABSTRACT

Vega has been shown to host multiple dust populations, including both hot exozodiacal dust at sub-au radii and a cold debris disc extending beyond 100 au. We use dynamical simulations to show how Vega's hot dust can be created by long-range gravitational scattering of planetesimals from its cold outer regions. Planetesimals are scattered progressively inwards by a system of 5–7 planets from 30 to 60 au to very close-in. In successful simulations, the outermost planets are typically Neptune mass. The back-reaction of planetesimals scattering causes these planets to migrate outwards and continually interact with fresh planetesimals, replenishing the source of scattered bodies. The most favourable cases for producing Vega's exozodi have negative radial mass gradients, with sub-Saturn- to Jupiter-mass inner planets at 5–10 au and outer planets of 2.5 – 20 M<sub> $\oplus$ </sub>. The mechanism fails if a Jupiter-sized planet exists beyond ~15 au because the planet preferentially ejects planetesimals before they can reach the inner system. Direct-imaging planet searches can therefore directly test this mechanism.

**Key words:** planets and satellites: detection – planets and satellites: dynamical evolution and stability – zodiacal dust – stars: Individual: Vega (alpha lyrae) – planetary systems – infrared: planetary systems.

## **1 INTRODUCTION**

Vega is the original rock star. Observations dating back 30 yr have shown a large excess in flux over the stellar photosphere at midinfrared and longer wavelengths (Aumann et al. 1984; Wilner et al. 2002; Sibthorpe et al. 2010). Additional excesses have been measured in the *K* band ( $\sim 2 \mu$ m) of  $\sim 1.3$  per cent the stellar flux (Absil et al. 2006; Defrère et al. 2011) and at 10–30 µm (excess of  $\sim 7$  per cent; Su et al. 2013).

The bright emission at  $\lambda \ge 50 \,\mu\text{m}$  comes from a resolved disc of more than 100 au in radius (Holland et al. 1998; Su et al. 2005). Cold dust is thought to be generated by the collisional grinding of a planetesimal belt located between roughly 60 and 120 au (Su et al. 2005; Müller, Löhne & Krivov 2010). The excess at 2  $\mu$ m is attributed to hot '-exozodiacal' dust at ~0.2 au that is dominated by small grains and cannot be in collisional equilibrium because of the too-short time-scales (Defrère et al. 2011). The warm excess at 10–30  $\mu$ m is spatially distinct from the hot emission and may represent an asteroid belt at ~14 au (Su et al. 2013), although the flux is likely contaminated by the tail of the hot dust (e.g. Lebreton et al. 2013) and, as we will explore below, perhaps by inward-scattered cold planetesimals.

For hot inner dust to come from the outer planetesimal belt, planets must continually interact with new material. A solution is for the outermost planet to migrate outwards. For a given range of planet masses and disc properties, outward migration is naturally maintained by the gravitational back-reaction of inward scattering of planetesimals (Gomes, Morbidelli & Levison 2004; Kirsh et al. 2009; Ormel, Ida & Tanaka 2012). Planets that are too low mass migrate rapidly but do not excite planetesimal eccentricities sufficiently to produce sufficient scattering. Planetesimal-driven

Here, we propose that Vega's hot exozodiacal dust may be supplied from its cold debris disc. Planetesimals are transported inwards by gravitational scattering by a system of planets. The outermost planet scatters planetesimals inwards and passes them off to the next planet until they reach the inner system; this is the same mechanism that delivers Jupiter-family comets to the inner Solar system (Levison & Duncan 1994) and is thought to be the main contributor to the Solar system's zodiacal dust (Nesvorný et al. 2010; Rowan-Robinson & May 2013). Planets on static orbits cannot sustain the inward-scattering rates needed to produce the observed exozodis because the outer planetesimal supply zones are depleted too rapidly (Bonsor, Augereau & Thébault 2012). The observed exozodis can also not be caused by dynamical instabilities among systems of giant planets (see Raymond, Armitage & Gorelick 2010) because the pulse of inward scattering is too short in duration (Bonsor, Raymond & Augereau 2013).

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migration naturally produces exozodis for a specific range of planet masses that depend on the disc's surface density (see Bonsor et al. 2014).

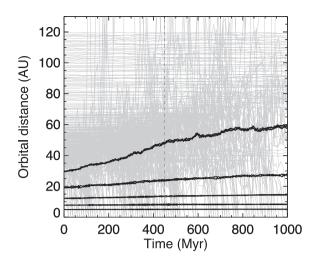
Our simulations constitute a proof-of-concept experiment. We determine the conditions needed for planetesimal-driven migration of a system of planets to generate a sufficient flux of inward-scattered planetesimals to roughly reproduce Vega's exozodi. We consider the outer debris disc as the sole source of planetesimals. We show that inward-scattered planetesimals can produce sufficient warm dust to explain the  $10-30 \,\mu\text{m}$  excess without invoking steady-state collisional grinding in a dynamically cold asteroid belt.

#### 2 SIMULATIONS

Our goal was to produce the observed hot exozodiacal dust. But the simulations also needed to match Vega's cold debris disc after 455 Myr of evolution. To simplify things, we focused on the bestfitting model of Müller et al. (2010) that reproduces the observations with a disc extending from 62 to 120 au containing 46.7  $M_{\oplus}$ , having ground down by ~20 per cent in mass over Vega's lifetime. The initial planetesimal discs were therefore somewhat more massive than the current one. Our setup was also designed to trigger outward planetesimal-driven migration of the outermost planet, which plays a role in sculpting the inner edge of the debris disc. We therefore started the planet significantly closer-in than the current planetesimal disc's inner edge.

Our fiducial setup had five planets placed between 5 and 30 au, evenly spaced in orbital period ratio  $(P_{i+1}/P_i \approx 1.95)$ . Our choice of five planets produced systems that would indeed trigger the migration-scattering mechanism, although our setup was far from optimal (see Bonsor et al. 2014, more compact orbital configurations are more effective at scattering). The number of planets is much less important than the fidelity of the chain of planets that passes planetesimals inwards. We varied the planets' masses between simulations as described below. The planets' initial orbits had randomly chosen eccentricities of less than 0.02 and randomly chosen inclinations of less than 1°. The planetesimal disc extended from 30 to 120 au and contained 90  $M_{\oplus}$  . It followed an  $r^{-1}$  surface density profile such that it started with  $60 M_{\oplus}$  within the 60–120 au belt of Müller et al. (2010). The disc contained 3000 planetesimals with randomly chosen eccentricities of up to 0.01 and inclinations of up to 0°.5. We adopted the Yoon et al. (2010) stellar parameters for Vega (age of 455  $\pm$  13 Myr and mass of 2.157  $\pm$  0.017 M<sub> $\odot$ </sub>).

To sustain an exozodi to its current age, the inner regions of the Vega system must be resupplied with material. In our simulations, we tracked the flux of planetesimals scattered into the inner regions, defined as interior to 3 au. Any particle that entered within this radius was removed from the simulation. The details regarding the manner in which scattered planetesimals may evolve collisionally, or due to radiative forces and/or sublimation, potentially leading to a population of small dust grains similar to those thought to produce the exozodi emission, is a complex problem that we do not address here. Instead we consider a lower limit on the mass of material that must be delivered to the inner regions, and favour simulations that transport more material inwards. Defrère et al. (2011) estimated Vega's exozodi to contain  $\sim 10^{-9} \, M_{\oplus}$  in dust with a characteristic lifetime of  $\sim 1$  yr. We adopt  $10^{-9} M_{\oplus} \text{ yr}^{-1}$  as a firm lower limit to the mass of material that must be delivered to the inner regions, and favour simulations that transport more material inwards. In some figures, a factor f represents the fraction of inward-scattered planetesimals that must be converted to hot dust



**Figure 1.** Evolution of the planets (periastron and apoastron in black) and planetesimals (semimajor axis in grey) in an example simulation. In radial order, the planets' masses are 50, 25, 25, 10, and  $10 M_{\oplus}$ . All planetesimals that entered within 3 au are shown but only 1 of every 30 other planetesimals is shown. Vega's estimated age is marked with a vertical dashed line.

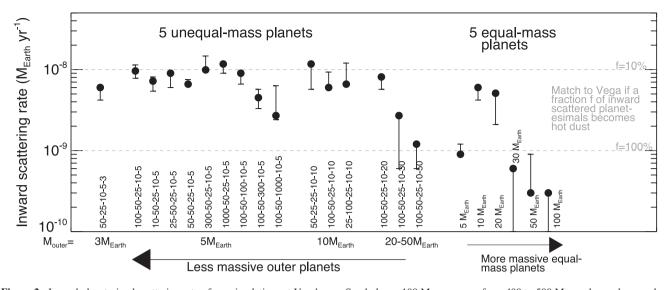
to match the observations (e.g., f = 10 per cent for a scattering rate of  $10^{-8} M_{\bigoplus} \text{ yr}^{-1}$ ).

Each simulation was integrated for 500 Myr to 1 Gyr using the hybrid MERCURY integrator (Chambers 1999). Planetesimals interacted gravitationally with the planets but not with each other. Collisions were treated as inelastic mergers. Particles were considered ejected and removed from the system if their radial distance exceeded  $10^4$  au. The timestep was set to 200 d, designed to conserve energy to better than one part in  $10^4$  outside 3 au (see appendix A in Raymond et al. 2011).

Fig. 1 shows the evolution of a simulation with planets of 50, 25, 25, 10, and  $10 \, M_{\oplus}$  (inner to outer). The outermost planet migrated from 30 au out to 59 au. The second- and third-outermost planets migrated from 19 to 28 au and from 12 to 14.5 au. The inner two planets did not migrate substantially although they did cross the 2:1 mean motion resonance at ~300 Myr, causing a modest increase in the planets' eccentricities. Due to their large Hill radii, none of the planet was hit by the most planetesimals but accreted just  $0.96 \, M_{\oplus}$ .

A total of  $7.6 \, M_{\oplus}$  was scattered inside 3 au in the simulation from Fig. 1 for an average of  $7.6 \times 10^{-9} \,\mathrm{M}_{\bigoplus} \,\mathrm{yr}^{-1}$ . Almost all of the in-scattered planetesimals originated between the portion of the planetesimal disc initially between 30 and 60 au. The outer planet migrated through this portion of the disc and removed more than 80 per cent of the mass, mostly by scattering planetesimals to larger orbital radii (including ejecting them) or by scattering them inwards towards the star. A low-density belt of planetesimals survived at 35-50 au. These were scattered by the outermost planet but not strongly enough to encounter the next planet (see Raymond et al. 2012). The disc is no longer as dynamically calm as the one from Müller et al. (2010), yet it retains similar properties. The median eccentricity of outer belt planetesimals is 0.05. The amount of mass from 60 to 120 au is  $64 M_{\oplus}$  without accounting for collisional grinding, accelerated by the larger planetesimal eccentricities (e.g. Wyatt 2008).

Fig. 2 shows the simulated inward-scattered flux at Vega's age. A first conclusion is that systems with equal-mass planets can only produce Vega's exozodi for planet masses of  $10-20 M_{\oplus}$ .



**Figure 2.** Inward planetesimal-scattering rate of our simulations at Vega's age. Symbols are 100 Myr averages from 400 to 500 Myr and error bars enclose the averages over the previous and subsequent 100 Myr intervals. Simulations with unequal-mass planets are sorted by the mass of the outermost planet and then by the mass of the next-outermost planet and are labelled with the masses of all five planets (inner-to-outer). The system from Fig. 1 is '50-25-25-10-10'. Simulations with equal-mass planets are labelled by the individual planet mass. The limits at  $10^{-9}[10^{-8}]M_{\oplus}$  yr<sup>-1</sup> can produce Vega's observed hot dust if 100 per cent [10 per cent] of the mass in inward-scattered planetesimals turns to hot dust.

Despite the low inward-scattering rates for equal-mass planets, they are roughly an order of magnitude higher than corresponding simulations with planets on static orbits from Bonsor et al. (2012). Planetesimal-driven migration is clearly a key mechanism in maintaining a steady inward flux of planetesimals. We note that planets evenly spaced in orbital period ratio are more efficient scatterers than planets separated by a fixed number of mutual Hill radii (Bonsor et al. 2014).

Many simulations with unequal-mass planets had inward-scattering rates large enough to reproduce Vega's exozodi (middle panel). The successful systems are broadly characterized by negative radial mass gradients. The inner planet is  $50 M_{\oplus}$  to (super-) Jupiter-mass, the outer 1–2 planets are  $5 - 20 M_{\oplus}$ , and the planets in between are  $10{-}50 M_{\oplus}$ . The exact ordering of planets is not important as long as this general trend is followed. The simulation from Fig. 1 has one of the highest inward-scattering rates, of  $1.2 \times 10^{-8} M_{\oplus} \text{ yr}^{-1}$ . The unequal-mass systems that did not generate large inward-scattering rates suffered from one of the two problems. Either the mass of their outermost planet was not between 5 and  $20 M_{\oplus}$  or the system contained a gas giant.

#### 2.1 Limits on the presence of a Jupiter-mass planet

We tested the effect of a Jupiter-mass planet on exozodi production. We started from a fiducial simulation with planet masses of 100, 50, 25, 10, and  $5 M_{\oplus}$  (inner to outer). The simulation had an inward-scattering rate of  $9.6 \times 10^{-9} M_{\oplus} \text{ yr}^{-1}$ . We performed five additional simulations. In each simulation, we increased the mass of one of the five planets to one Jupiter-mass. Fig. 3 shows the inward-scattering rate in these systems. Increasing the mass of the innermost or second-innermost planet did not have a large effect on the scattering rate. For a more distant Jupiter, the rate of inward-scattered planetesimals was significantly depressed. The rate was cut in half with a Jupiter in the middle of the system (at 12.2 au). The rate drops is borderline for a Jupiter at 19.2 au and below the critical limit at 30 au.

Inward scattering rate (10<sup>-9</sup> M<sub>Earth</sub> yr<sup>-1</sup> 14 12 10 8 6 4 2 f=100% 0 5 35 10 15 30 0 20 25 Orbital radius of Jupiter-mass planet (AU)

**Figure 3.** Inward-scattering rate at Vega's age for systems in which each of the five planets in a fiducial simulation was successively replaced with a Jupiter-mass planet. The symbols are averages from 400–500 Myr and the error bars encompass rates over the two previous and successive 100 Myr intervals. The horizontal lines mark the rates for which 10 per cent and 100 per cent of in-scattered planetesimals must be converted to hot dust to match the observations.

With its high escape speed  $(60 \text{ km s}^{-1})$ , a Jupiter-mass planet ejects planetesimals more efficiently than it passes them inwards. It disrupts the supply chain between the outer and inner planetary systems. The mechanism of inward scattering is thwarted by a gas giant planet beyond ~15 au. This is testable: no gas giant planet can exist beyond ~15 au for our mechanism to operate. Current observations constrain the mass of a planet to be less than 3–10 $M_J$ beyond ~10 au (Itoh, Oasa & Fukagawa 2006; Marois et al. 2006; Heinze et al. 2008). New direct-imaging instruments are coming online soon, specifically SPHERE (Beuzit et al. 2008) and Gemini Planet Imager (Macintosh et al. 2008). Observations may directly constrain our model in the coming years.

#### 2.2 Testing different planetesimal belt configurations

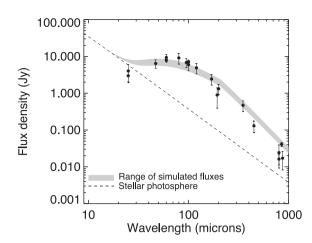
To show that the mechanism is robust we ran two additional sets of simulations with different planetesimal disc properties. The first set assumed a lower mass disc with  $50 M_{\oplus}$  initially contained between 30 and 120 au. The second set assumed a lower mass and colder disc containing  $50 M_{\oplus}$  between 45 and 120 au. Each disc contained 3000 planetesimals and the same orbital distribution as in the fiducial disc. For the lower mass disc, we ran simulations with planets with similar mass gradients to the most successful simulations in the fiducial disc as well as equal-mass planets with  $M_p = 2.5-40 M_{\oplus}$ . For the cold disc, we kept the innermost planet at 5 au while maintaining dynamical contact with the outer disc. We placed the planets between 5 and 45 au, again evenly spaced in orbital period ratio. In some simulations, we included five planets with  $P_{i+1}/P_i = 2.3$  and in others seven planets with  $P_{i+1}/P_i = 1.7$ .

The simulations in cold discs produced inward-scattering rates as high as  $7 \times 10^{-9} \, M_{\oplus} \, yr^{-1}$ . The most successful simulations had seven planets including a Jupiter–Saturn-like inner pair and outer planets with masses between 2.5 and  $10 \, M_{\oplus}$ . Although fiveplanet systems generally had much lower inward-scattering rates, the most successful five-planet system – with planet masses of 300, 100, 50, 25, and  $10 \, M_{\oplus}$  from innermost to outermost – had a comparable scattering rate to the successful seven-planet systems. In the low-mass disc, only two simulations had high enough inwardscattering rates to match the  $10^{-9} \, M_{\oplus} \, yr^{-1}$  needed to generate Vega's exozodi. Both successful simulations had negative radial mass gradients, with (inner to outer) planets masses of 300, 100, 50, 25,  $10 \, M_{\oplus}$ , and 100, 50, 25,  $10, 5 \, M_{\oplus}$ .

#### 2.3 Matching observations

We now estimate the effect that scattered planetesimals have on the observations. Transforming a planetesimal population into a dust distribution is a complex process (e.g. Dominik & Decin 2003; Wyatt 2008; Krivov 2010). We use a simplified approach. Our goal is not to perfectly reproduce the observations but to verify that observations of our simulated planetary systems would not differ enormously from those of Vega. We used the code from Raymond et al. (2011, 2012), which follows the procedure of Booth et al. (2009), based on the model of Wyatt et al. (2007).

Each simulated planetesimal was assumed to represent a distribution of objects following a collisional size-frequency with  $dN/dD \propto D^{-3.5}$  (Dohnanyi 1969). Each planetesimal is heated by the star and cools as a blackbody. We took into account Vega's stellar properties (Yoon et al. 2010) to derive an approximate blowout size of  $D_{\rm bl} \approx 0.8 \left( L_{\star}/L_{\odot} \right) \left( M_{\star}/M_{\odot} \right)^{-1} \approx 15 \,\mu {\rm m}.$ The surface area per unit mass of our collisional distribution is  $\sigma/M = 0.07 \,\mathrm{au^2}\,\mathrm{M_{\bigoplus}}^{-1}$  (assuming  $D_{\mathrm{max}} = 2000$  km). The collisional time-scale  $t_{coll}$  of the largest bodies was calculated and the planetesimal mass was decreased in time as  $M_{\rm p}(t) \propto \frac{1}{1+t/t_{\rm coll}}$ . This is not strictly consistent with the dynamical integration in which the planetesimal mass did not change, but the cumulative decrease in  $M_p$  is relatively small (typically 10–20 per cent). We assumed a fixed value of  $Q_D^{\star} = 200 \,\mathrm{J\,kg^{-1}}$  (Wyatt et al. 2007). This does not take into account realistic dust emission, so fluxes are to be trusted within an order of magnitude (Booth et al. 2009).



**Figure 4.** Vega's observed SED (data compiled from table 1 in Müller et al. 2010) compared with SEDs generated by our simulations. The shaded region encloses the fluxes generated by three simulations that produced large inward-scattering fluxes: one in the fiducial disc (the upper bound), one in the cold disc, and one in the low-mass disc.

Fig. 4 shows that the spectral energy distribution (SED) of our simulations provides a decent match to Vega's dust flux in the midinfrared. There is a small excess in the fiducial simulations (the upper bound of the shaded area in Fig. 4) at 25 and 47 µm. The excess comes from planetesimals scattered inwards on Centaur-like orbits. These planetesimals are hotter than those in the outer disc and contaminate the SED starting at slightly shorter wavelength than the observed infrared excess. They also generate warm dust at a similar level to the observations. Inward-scattered planetesimals contribute a few to 10 per cent of the stellar flux at  $10-30 \,\mu m$ , the same order of magnitude as the Su et al. (2013) claimed asteroid belt. In-scattered planetesimals may account for both Vega's warm and hot dust populations. However, we caution that the true dust production of these planetesimals is highly uncertain. As they may not be in collisional equilibrium as assumed, inward-scattered planetesimals may produce less dust than in our simple calculations. In addition, any terrestrial planets inside 3 au would affect the long-term dynamics of both inward-scattered planetesimals and in-spiralling dust.

The dust SEDs from the successful simulations in cold and lowmass discs somewhat underestimated the flux at  $50-100 \mu m$  because the planetesimal mass was somewhat below the expected one (the lower bound of the shaded region in Fig. 4).

#### **3 DISCUSSION**

We have shown that Vega's exozodi can be explained by a system of planets with orbital radii of roughly 5–60 au. Planets scatter planetesimals from the outer planetesimal disc into the inner planetary system. The outermost planets undergo planetesimal-driven migration outwards into the disc, continually encounter fresh material, and can sustain a substantial inward-scattering rate for the age of the system and beyond (see Bonsor et al. 2014, for details of this process).

For the mechanism to operate, the system must meet the following requirements. First, the outermost 1–2 planets must have masses such that planetesimal-driven migration is triggered and sustained. Successful simulations had  $M_{outer} = 2.5-20 \,\mathrm{M_{\bigoplus}}$ . The inner planets' masses have only a minor effect, although the highest scattering rates were for simulations with higher mass (Saturn- to Jupiter-mass) planets closer-in. Secondly, while the inner system may host a gas giant, an ~Jupiter-mass planet cannot orbit beyond  $\sim$ 15 au as it too efficiently ejects planetesimals from the system and breaks the inward-scattering chain (see Fig. 3). This prediction does not violate current constraints (Itoh et al. 2006; Marois et al. 2006; Heinze et al. 2008) and may be tested with upcoming instruments. Thirdly, our model requires that a substantial fraction of the mass in planetesimals scattered within 3 au must end up as fine-grained dust. This fraction is  $\sim 10-30$  per cent for the simulations we have deemed successful. If we take our simplified dust flux calculation at face value, our simulations provide a modestly good match to the observed SED (see Fig. 4). Although our simulations somewhat overestimate the flux at 25-50 µm, dust produced by in-scattered planetesimals may offer an alternate explanation for the 10-30 µm excess detected by Su et al. (2013). Planetesimal-driven migration of a system of 5-7 planets in a lower mass disc and colder disc can also yield reasonable inward-scattering rates and SEDs.

## ACKNOWLEDGEMENTS

We thank M. Booth and J.-C. Augereau for interesting discussions. AB acknowledges the support of the ANR-2010 BLAN-0505- 01 (EXOZODI) and funding from NERC. This work was performed as part of the NASA Astrobiology Institute's Virtual Planetary Laboratory Lead Team, supported by the NASA under Cooperative Agreement no. NNA13AA93A. This Letter was entirely written while SNR was standing and is therefore a 'sit-free' Letter.

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